

Uncertainty Quantification for Reaction Theory



A.E. Lovell
April 21, 2020
Reaction Seminar



Managed by Triad National Security, LLC for the U.S. Department of Energy's NNSA

Outline

- Sources of uncertainty in few-body reactions
 - Uncertainty quantification methods
 - χ^2 minimization
 - Bayesian methods
 - Results
 - χ^2 minimization – $^{12}\text{C}(d,p)$
 - Bayesian methods – $^{48}\text{Ca}(d,p)$
 - Comparison between methods
 - Connections and outlook
- A.E. Lovell, F.M. Nunes, *J. Phys. G* **42**, 034014 (2015)
 - A.E. Lovell, F.M. Nunes, J. Sarich, S. M. Wild, *PRC* **95**, 024611 (2017)
 - A. E. Lovell, F. M. Nunes, *PRC* **97**, 064162 (2018)
 - G. B. King, A. E. Lovell, F. M. Nunes, *PRC* **98**, 044628 (2018)
 - G. B. King, A. E. Lovell, L. Neufcourt, F. M. Nunes, *PRL* **122**, 232501 (2019)
 - M. Catacora-Rios, G. B. King, A. E. Lovell, F. M. Nunes, *PRC* **100**, 064615 (2019)
 - Lovell, Nunes, Catacora-Rios, King, Carriello, *in prep*
 - Lovell, Mohan, Talou, *in prep*

Uncertainty quantification in nuclear theory

Enhancing the interaction between nuclear experiment and theory through information and statistics (ISNET) 2015

[2015 J. Phys. G: Nucl. Part. Phys. 42 Issue 3](#) Guest Editors: David Ireland and Witek Nazarewicz

Focus on further enhancing the interaction between nuclear experiment and theory through information and statistics (ISNET 2.0)

Now

Guest Editors

Dick Furnstahl Ohio State University

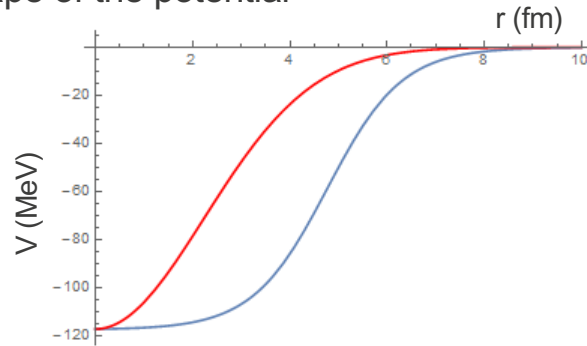
David Ireland University of Glasgow

Daniel Phillips Ohio University

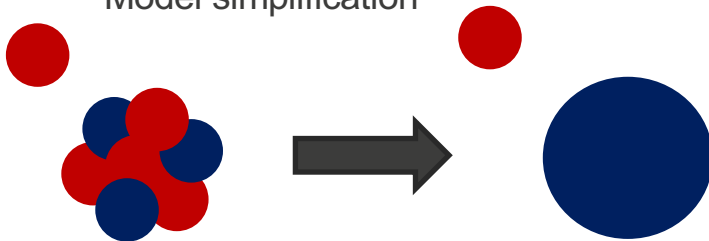
Types of uncertainty in few-body reaction theories

Systematic Uncertainties

Shape of the potential

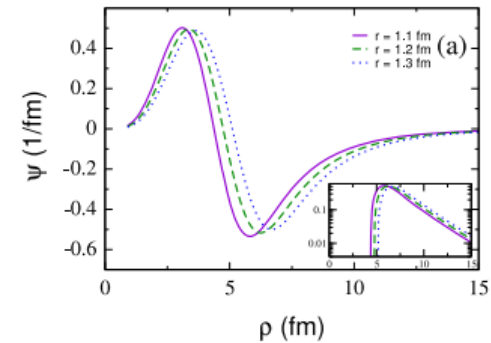


Model simplification



Statistical Uncertainties

Constraints on parameters



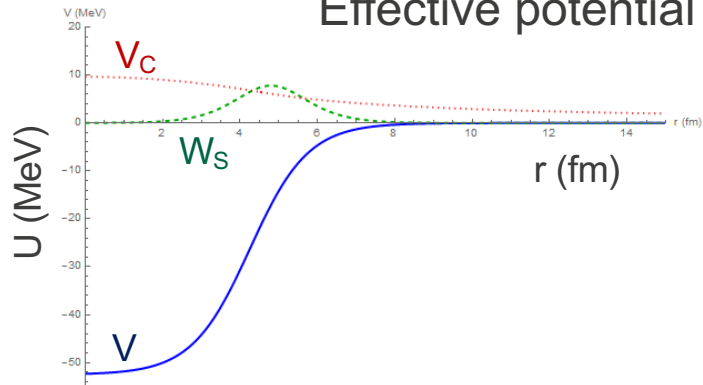
Convergence of functions

$$\sum_{k=1}^{k_{max}} f_k(x) \approx \sum_{k=1}^{\infty} f_k(x)$$

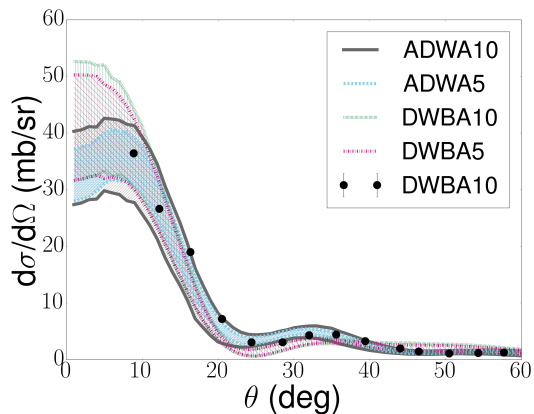
A.E. Lovell and F.M. Nunes J. Phys. G **42** 034014 (2015)

Sources of uncertainty in few-body reaction theories

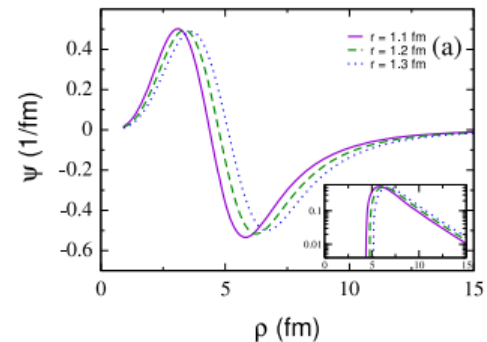
Effective potential



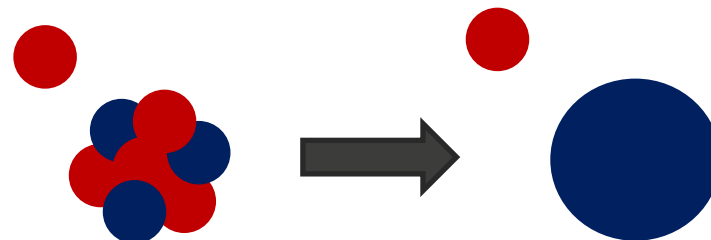
Few-body approximations made



Structure functions

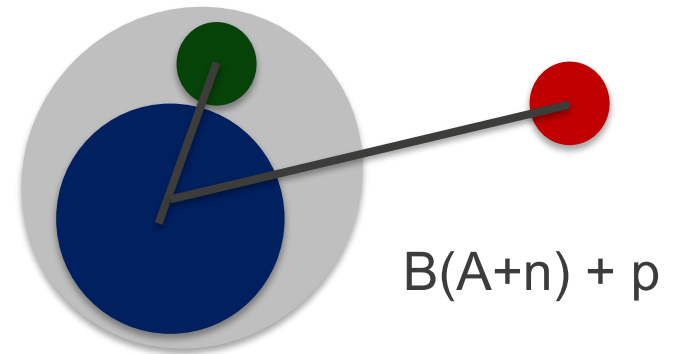
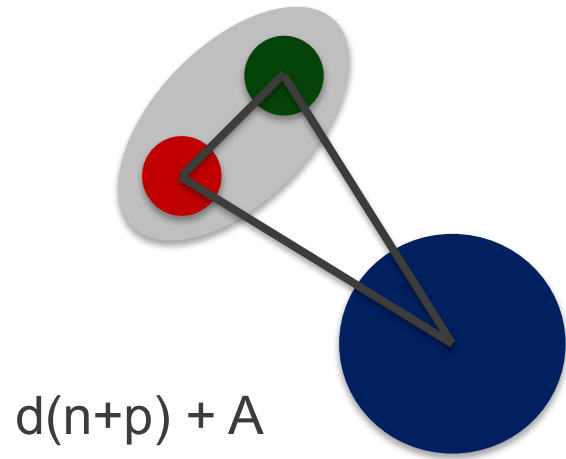


Missing degrees of freedom



A.E. Lovell and F.M. Nunes J. Phys. G **42** 034014 (2015)

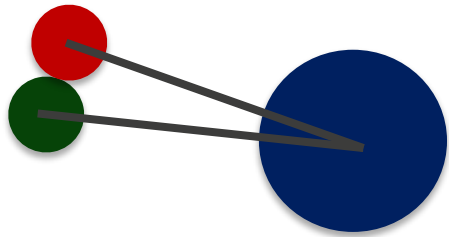
Reactions of interest: $A(d,p)B$



Few-body models for single-nucleon transfer reactions

$$[T_R + H_{\text{int}}(\vec{r}) + V_{pA} + V_{nA} - E]\Psi^{\text{exact}}(\vec{r}, \vec{R}) = 0 \quad H_{\text{int}} = T_r + V_{np}(r)$$

Adiabatic Wave Approximation (ADWA)



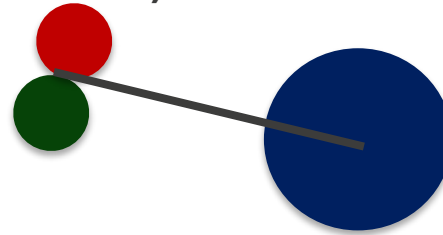
Nucleon-target potentials take into account deuteron break-up

$$[T_R + \epsilon_0 + V_{pA} + V_{nA} - E]\Psi^{\text{ad}}(\vec{r}, \vec{R}) = 0.$$

$$\Psi^{\text{ad}}(\vec{r}, \vec{R}) = \phi_0(\vec{r})\chi_0^{\text{ad}}(\vec{R}) + \sum_{i>0} \phi_i(\vec{r})\chi_i^{\text{ad}}(\vec{R})$$

Finite-range, Johnson and Tandy, *Nucl. Phys. A* **235** 56 (1974)

Distorted-Wave Born Approximation (DWBA)



Uses an effective deuteron-target potential

$$\Psi^{\text{exact}}(\vec{r}, \vec{R})$$



$$\Phi_{I_b I_a}(\vec{r}_i)\chi_i(\vec{R}_i)$$

np bound state

Deuteron scattering state

I.J. Thompson and F.M. Nunes, *Nuclear Reactions for Astrophysics*, (Cambridge University Press, Cambridge, 2009)

Phenomenological optical model potentials

$$U(r) = V(r) + iW(r) + (V_{so}(r) + iW_{so}(r))(\mathbf{l} \cdot \mathbf{s}) + V_C(r)$$

Volume Term

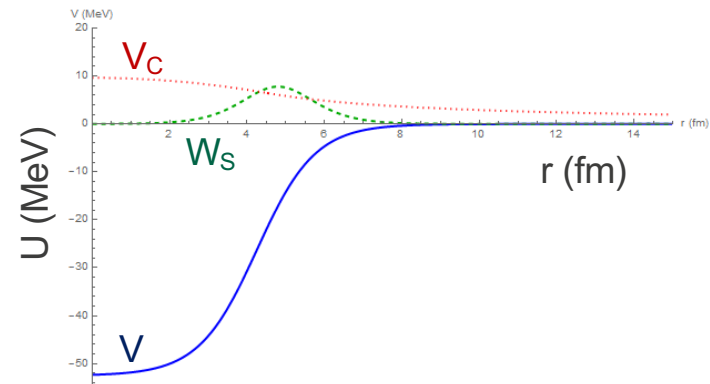
$$V(r) = f(r; V_o, R_o, a_o)$$

$$f(r; V_o, R_o, a_o) = -\frac{V_o}{1 + e^{(r-R_o)/a_o}}$$

≈ 6-12 free parameters

Surface and Spin-Orbit Terms

$$V(r) = \frac{d}{dr} f(r; V_o, R_o, a_o)$$



I.J. Thompson and F.M. Nunes, *Nuclear Reactions for Astrophysics*, (Cambridge University Press, Cambridge, 2009)

χ^2 minimization and covariance propagation

- In fitting the model to the observables, the goal is to minimize the residuals

$$\chi^2 = \sum_{i=1}^n \left(\frac{m(\mathbf{x}; \theta_i) - d_i}{\sigma_i} \right)^2 \quad [m(\hat{\mathbf{x}}; \theta_1) - d_1, \dots, m(\mathbf{x}; \theta_n) - d_n]^T \sim \mathcal{N}(0, \Sigma)$$

$\hat{\mathbf{x}} \longrightarrow$ Best fit set of parameters

- Assume that the parameters are normally distributed around the best-fit set, pull parameter sets and run them through the model

$$\mathcal{N}(\hat{\mathbf{x}}, \mathbb{C}_p) = \frac{1}{\sqrt{2\pi|\mathbb{C}_p|}} e^{-\frac{1}{2}(\mathbf{x}-\hat{\mathbf{x}})^T \mathbb{C}_p^{-1}(\mathbf{x}-\hat{\mathbf{x}})} \quad s^2 = \frac{1}{n-p} \sum_{i=1}^n \left(\frac{m(\mathbf{x}; \theta_i) - d_i}{\sigma_i} \right)^2 \quad \mathbb{C}_p \rightarrow s^2 \mathbb{C}_p$$

- 95% confidence intervals are constructed by removing the top 2.5% and bottom 2.5% of the calculations at each angle

Makes use of I. Thompson's fresco/sfresco:
www.fresco.org.uk

Bayesian optimization

Bayes' Theorem

$$P(\mathcal{H}|\mathcal{D}) = \frac{P(\mathcal{D}|\mathcal{H})P(\mathcal{H})}{P(\mathcal{D})}$$

Posterior – probability that the model/parameters are correct after seeing the data

Prior – what is known about the model/parameters before seeing the data

Likelihood – how well the model/parameters describe the data

Evidence – marginal distribution of the data given the likelihood and the prior

Parameter sets are accepted through the Markov Chain Monte Carlo algorithm (MCMC)

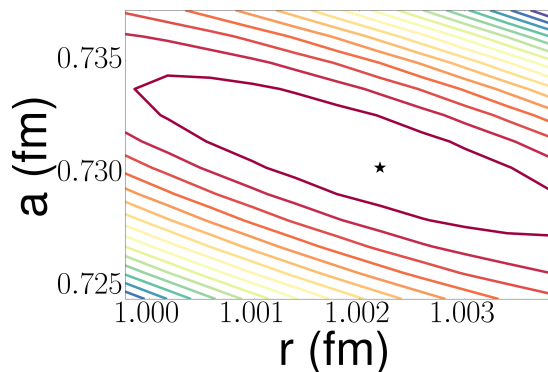
Confidence intervals are defined by the densest region of each parameter distribution or observable calculation

$$p(H) \propto \prod_{i=1}^{N_p} \exp \left[-\frac{(x^i - x_0^i)^2}{(x_0^i)^2} \right]$$

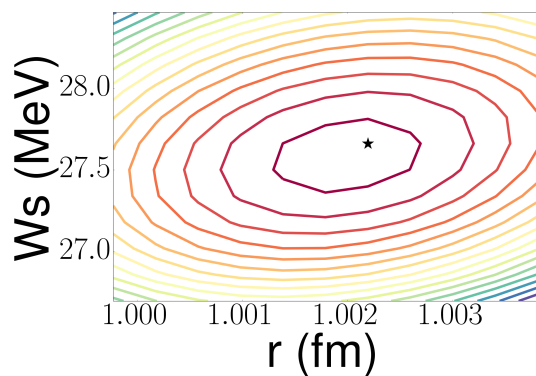
Built on top of I. Thompson's fresco:
www.fresco.org.uk

Fitting elastic scattering with χ^2 minimization

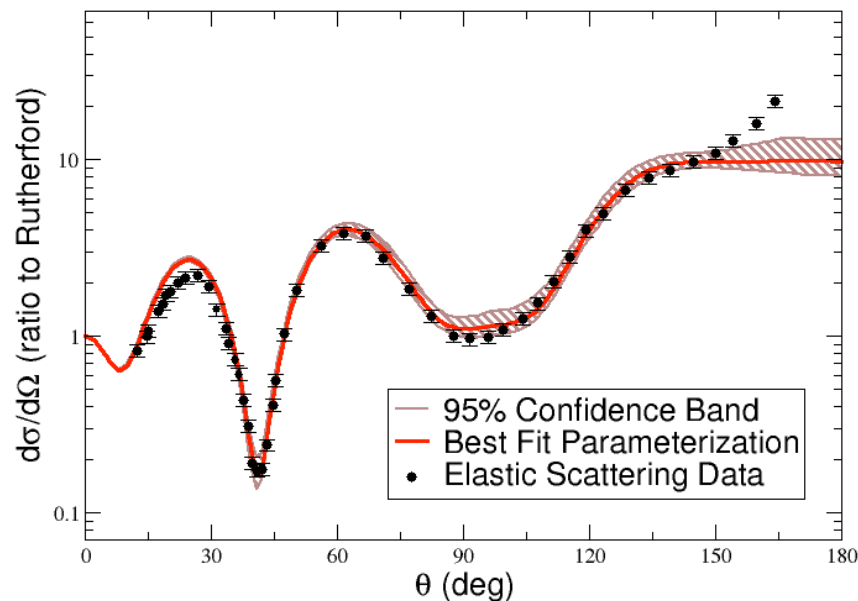
Correlated parameters



Uncorrelated parameters



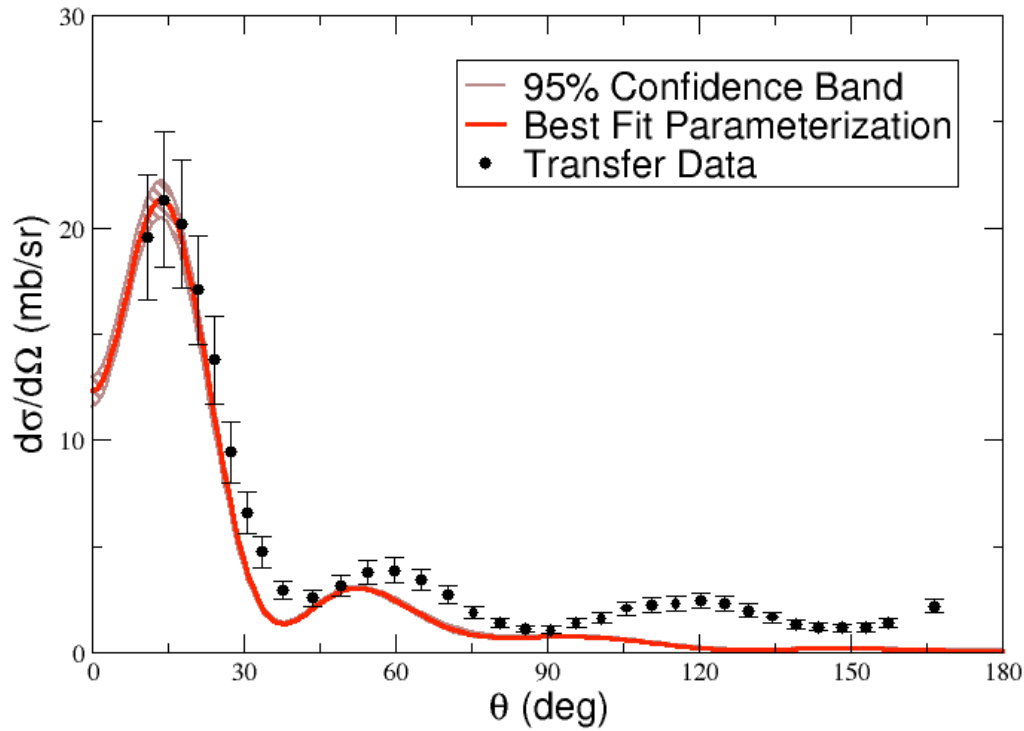
$^{12}\text{C}(d,d)$ @ 11.8 MeV



A.E. Lovell, F.M. Nunes, J. Sarich, S. M. Wild, *PRC* **95**, 024611 (2017)

Propagating to the transfer cross section (DWBA)

$^{12}\text{C}(d,p)^{13}\text{C}(\text{g.s.}) @ 11.8 \text{ MeV}$

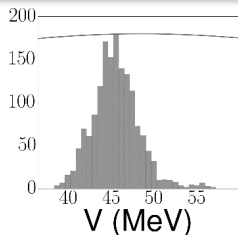


A.E. Lovell, F.M. Nunes, J. Sarich, S. M. Wild, *PRC* **95**, 024611 (2017)

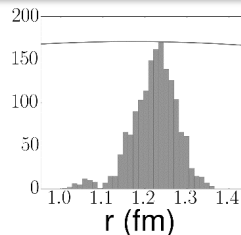
Bayesian posterior distributions

$^{48}\text{Ca}(n,n)$ @ 12 MeV

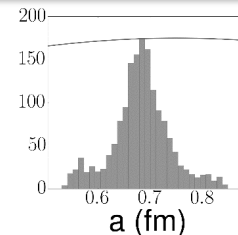
Real
Volume



$\mu=45.51$
 $\sigma=2.74$
(MeV)

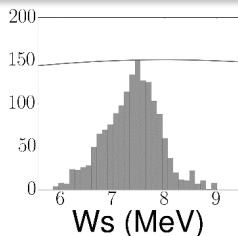


1.22
0.05
(fm)

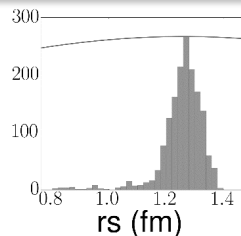


0.68
0.06
(fm)

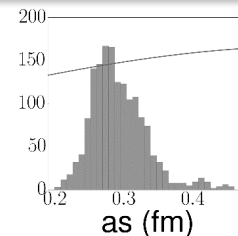
Imaginary
Surface



7.38
0.54
(MeV)

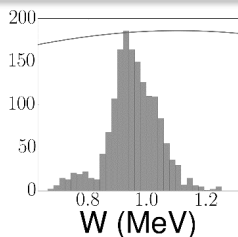


1.25
0.08
(fm)

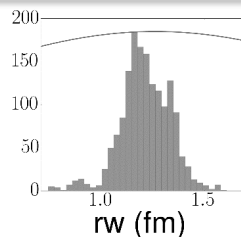


0.29
0.04
(fm)

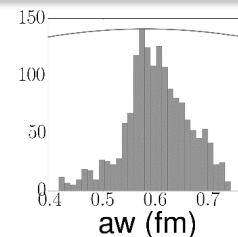
Imaginary
Volume



0.95
0.09
(MeV)



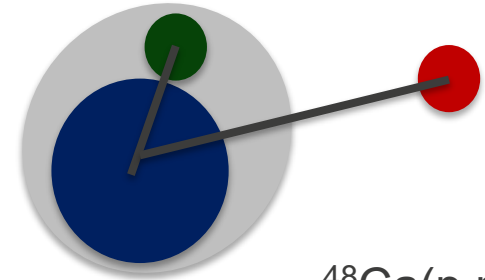
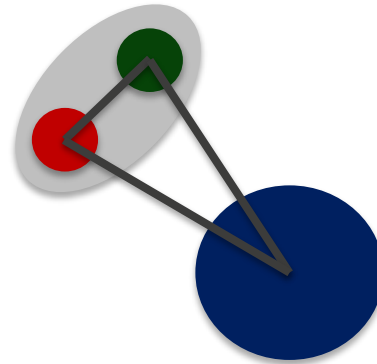
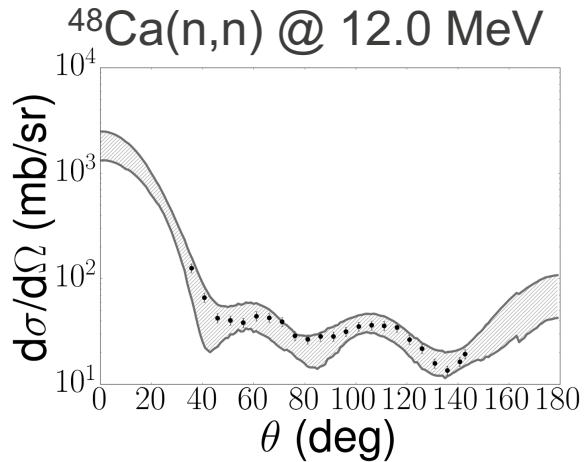
1.21
0.12
(fm)



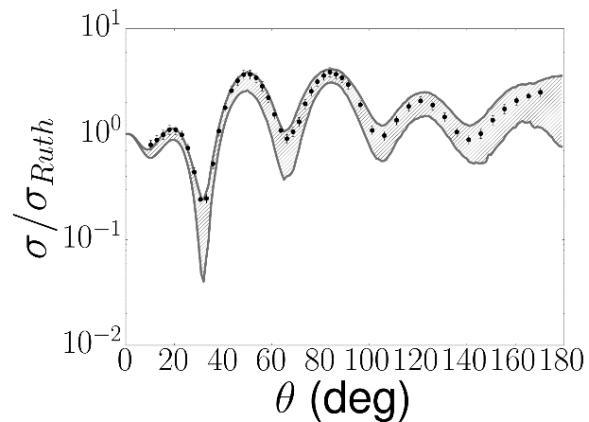
0.60
0.06
(fm)

A. E. Lovell, F. M. Nunes, *PRC* **97**, 064162 (2018)

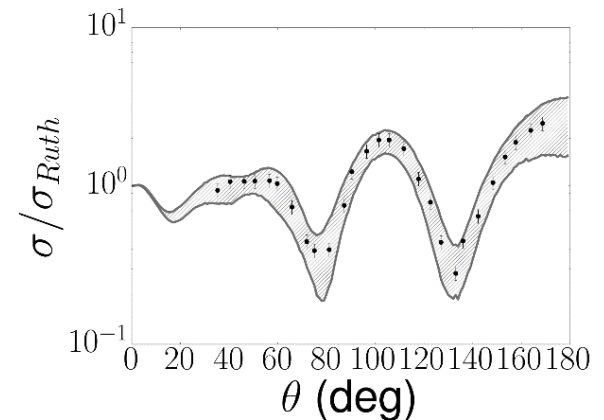
Bayesian elastic scattering confidence intervals



$^{48}\text{Ca}(p,p)$ @
25.0 MeV



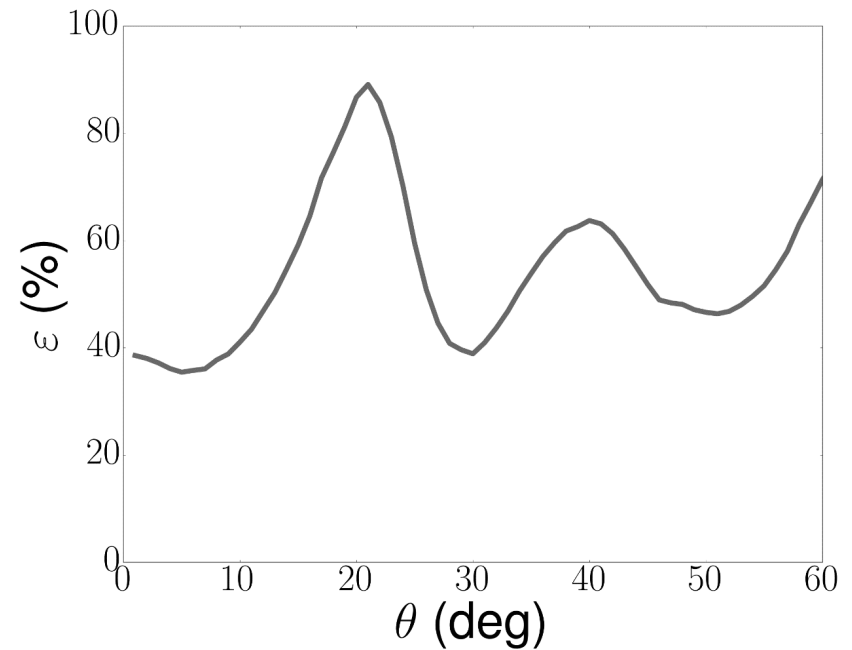
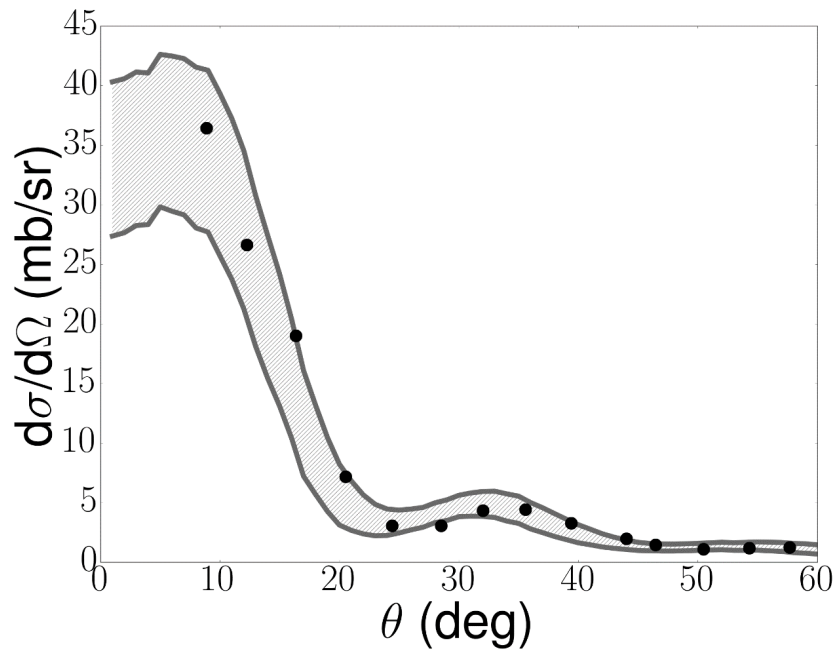
$^{48}\text{Ca}(p,p)$ @
14.03 MeV



A. E. Lovell, F. M. Nunes,
PRC **97**, 064162 (2018)

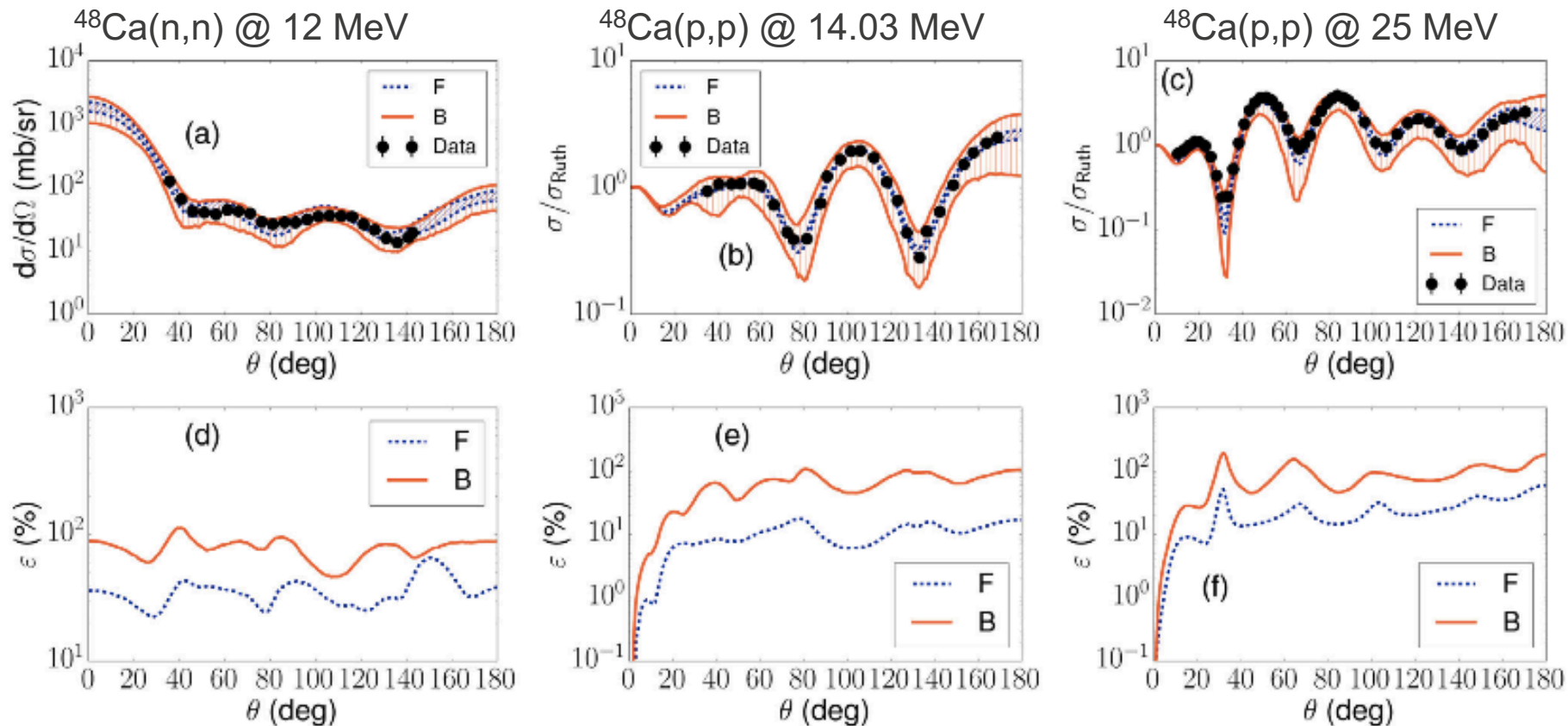
Propagation to the transfer cross section (ADWA)

$^{48}\text{Ca}(d,p)^{49}\text{Ca}(\text{g.s.}) @ 24 \text{ MeV}$



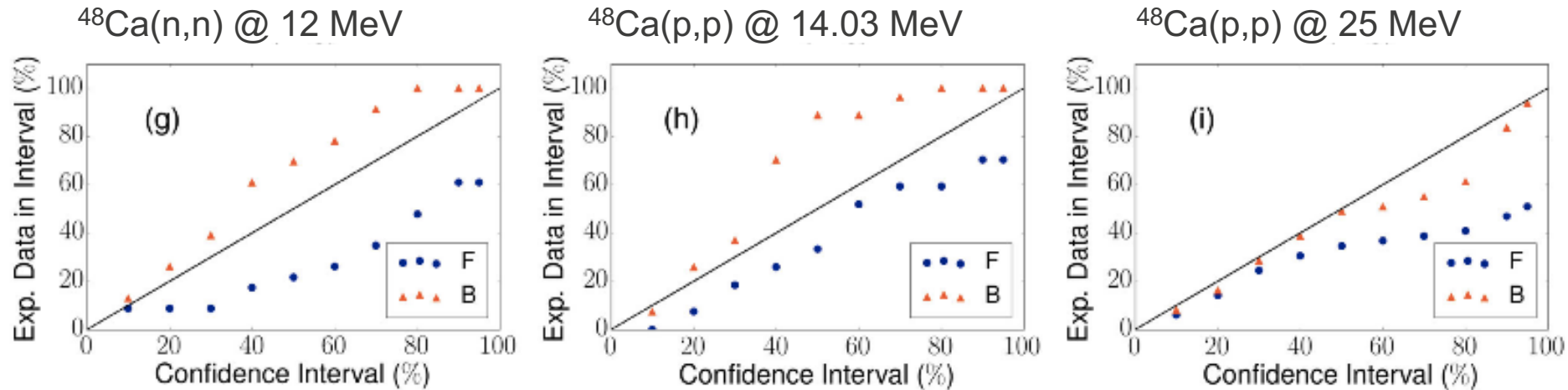
A. E. Lovell, F. M. Nunes, *PRC* **97**, 064162 (2018)

Direct comparison between χ^2 and Bayesian: cross section uncertainties



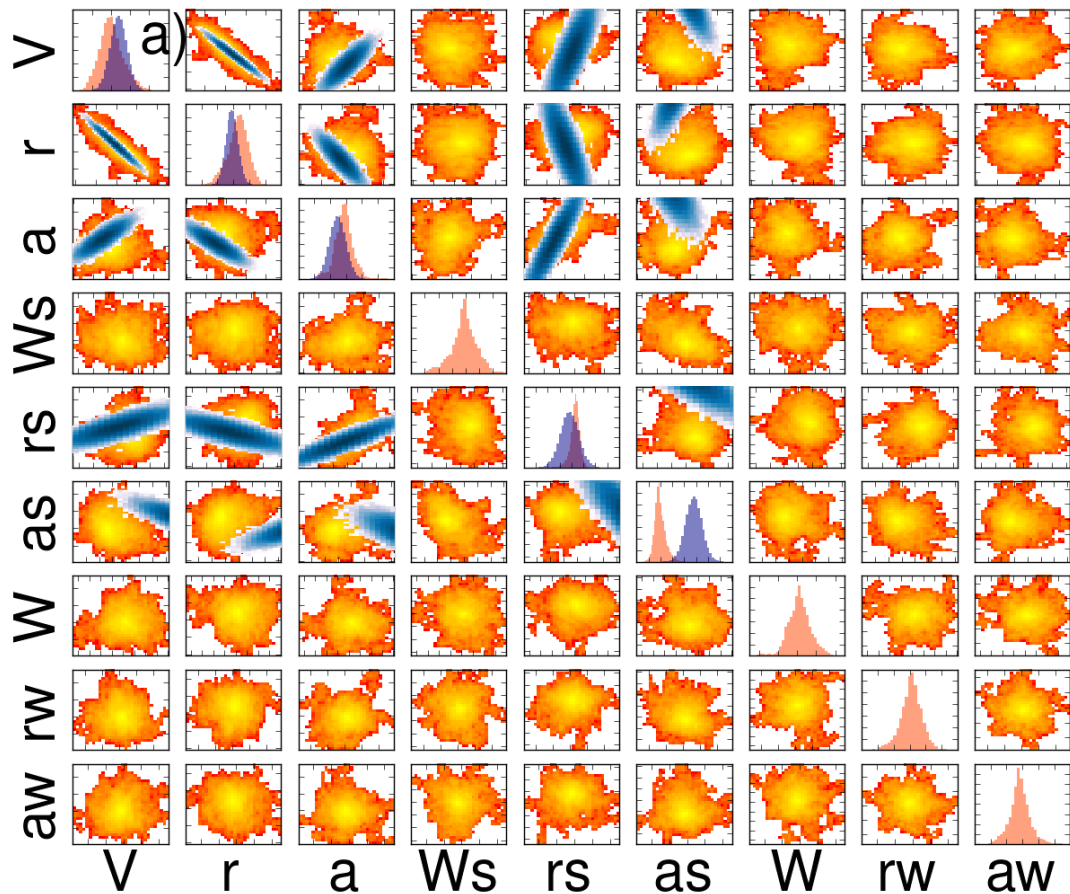
G.B. King, A.E. Lovell, L. Neufcourt, F.M. Nunes, *PRL* **122**, 232502 (2019)

Direct comparison between χ^2 and Bayesian: confidence in uncertainties



G.B. King, A.E. Lovell, L. Neufcourt, F.M. Nunes, *PRL* **122**, 232502 (2019)

Direct comparison between χ^2 and Bayesian: parameter correlations



$^{48}\text{Ca}(n,n)^{48}\text{Ca}$
@ 12 MeV

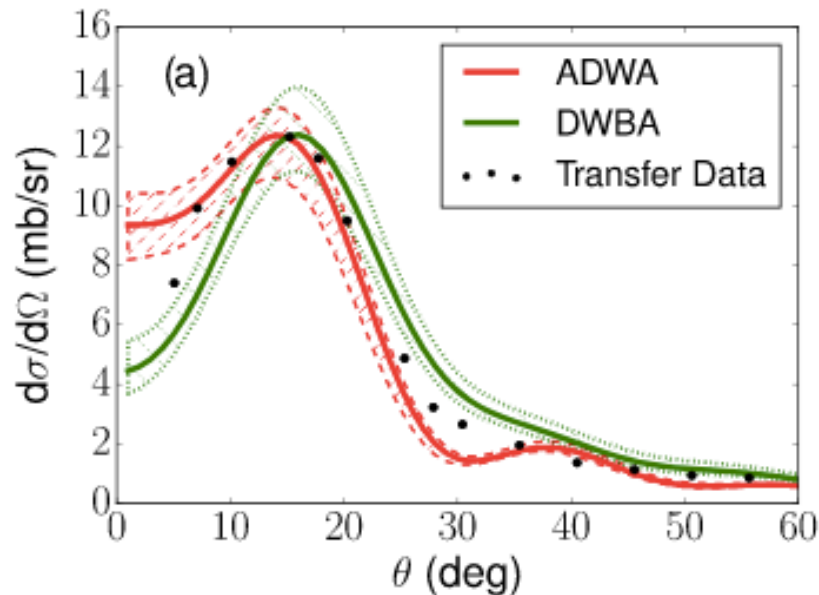
χ^2
minimization

Bayesian

G.B. King, A.E. Lovell, L. Neufcourt, F.M. Nunes, *PRL* **122**, 232502 (2019)

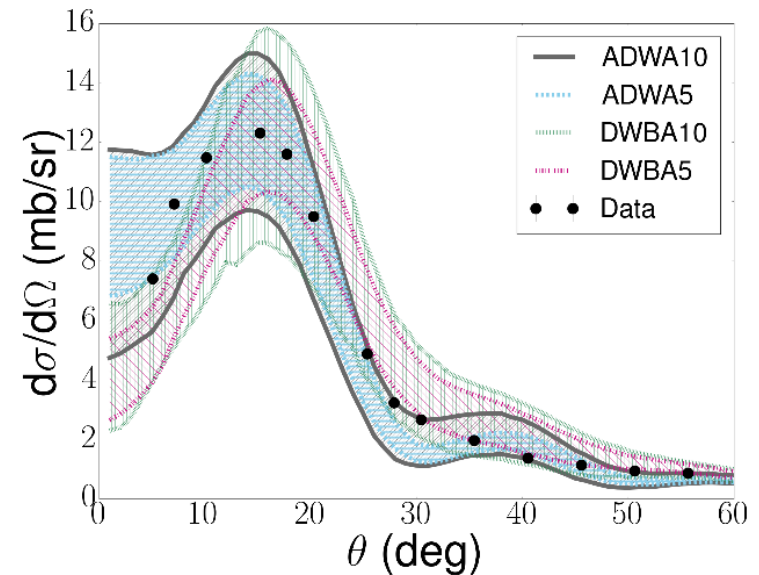
Other sources of uncertainty: few-body approximations

$^{90}\text{Zr}(d,p)^{91}\text{Zr}(\text{g.s.})$ @ 22.7 MeV
 χ^2 minimization



G.B. King, A.E. Lovell, F.M. Nunes,
PRC **98**, 044623 (2018)

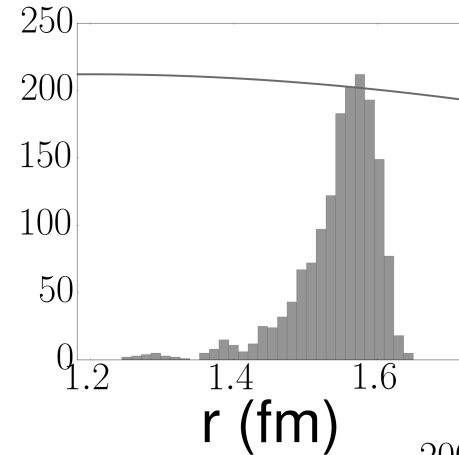
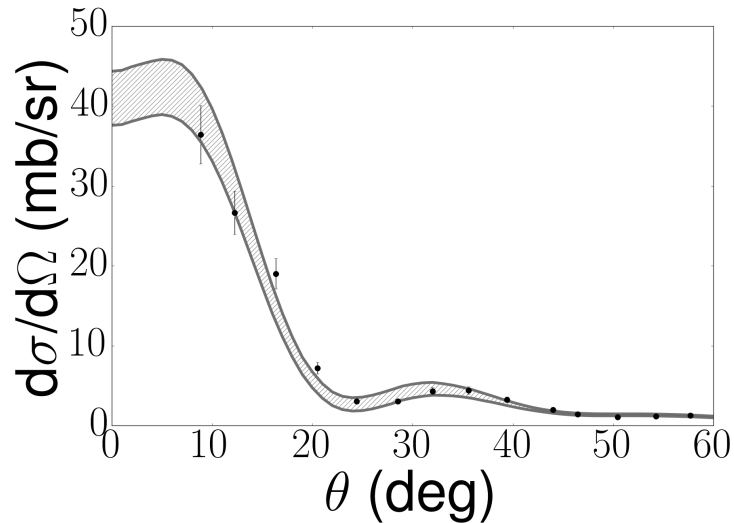
$^{90}\text{Zr}(d,p)^{91}\text{Zr}(\text{g.s.})$ @ 22.0 MeV
Bayesian



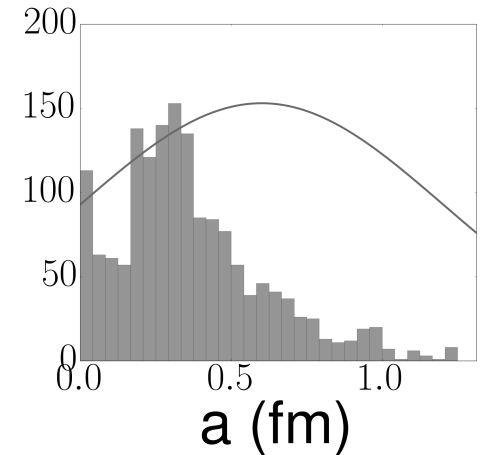
A.E. Lovell, MSU/NSCL Thesis (2018)

Other sources of uncertainty: structure functions

$^{48}\text{Ca}(d,p)^{49}\text{Ca}(\text{g.s.}) @ 19.3 \text{ MeV}$



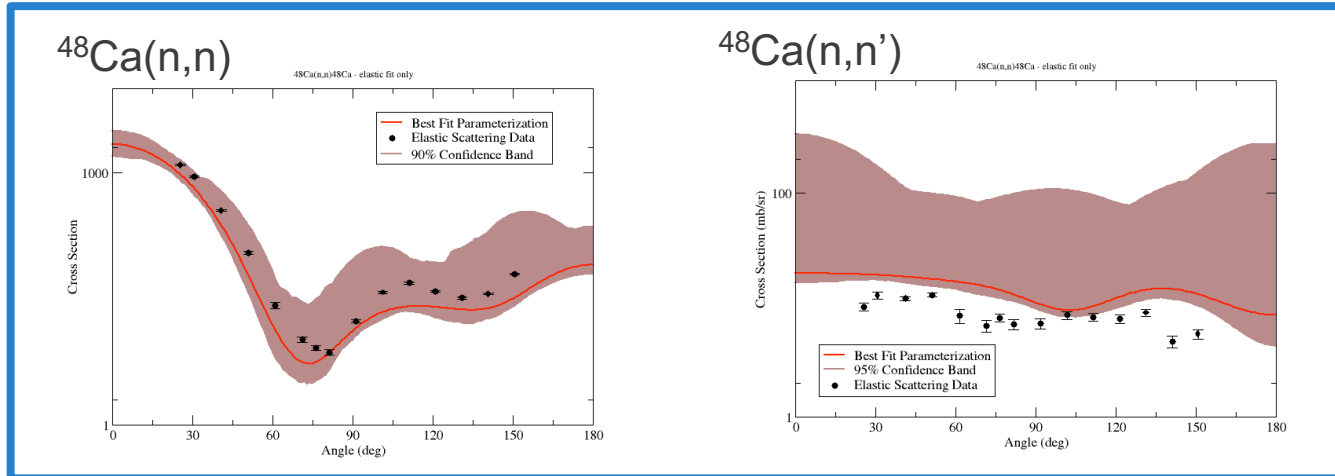
Preliminary work
to constrain the
single-particle
geometry



Much better to use an
independent constraint, such as
the asymptotic normalization
coefficient (ANC)

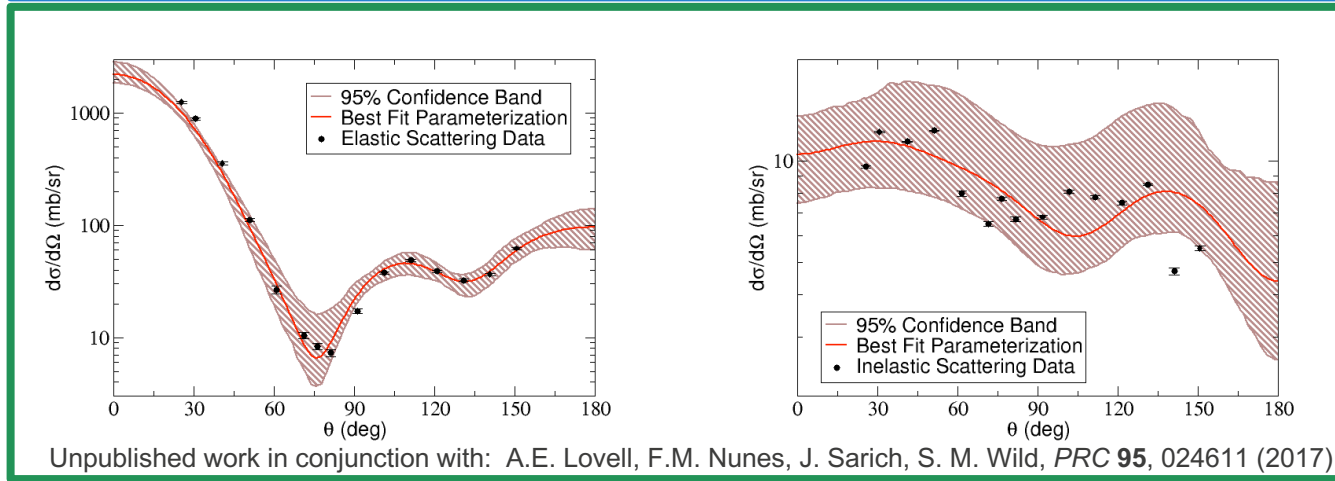
A.E. Lovell, unpublished
In progress work of N. Carriello

Other sources of uncertainty: degrees of freedom left out of the model



χ^2 minimization

Only elastic scattering fitted



Elastic and inelastic fitted with coupled channel calculation

Unpublished work in conjunction with: A.E. Lovell, F.M. Nunes, J. Sarich, S. M. Wild, *PRC* **95**, 024611 (2017)

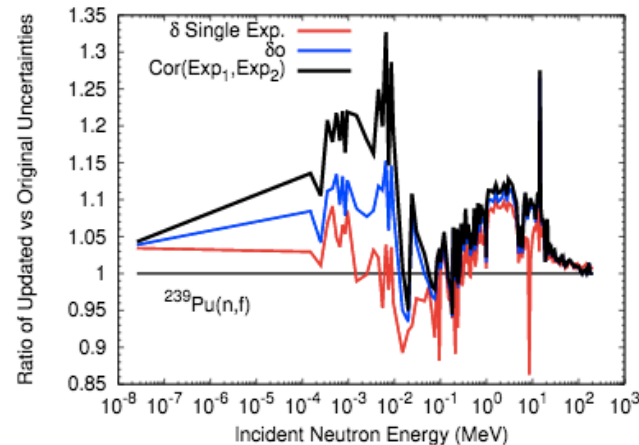
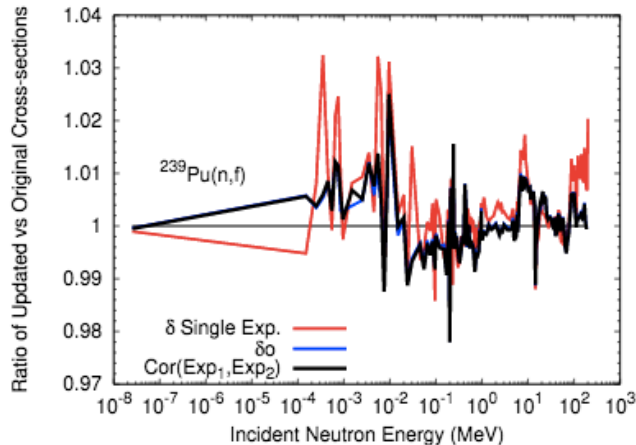
Nuclear data evaluations at LANL

ENDF/B-VIII.0 ENDF/B-VIII.0 released February 2, 2018

ENDF/B-VIII.0 fully incorporates the new Neutron Data Standards, includes improved neutron scattering data and uses new evaluated data from the Coordinated International Evaluation Library Organization (CIELO) pilot project for neutron reactions on ^{235}U , ^{238}U and ^{239}Pu . Notable advances include updated evaluated data for structural materials, actinides, fission energy release, prompt fission neutron spectra, thermal neutron scattering data, and charged-particle reactions.

D.A. Brown, *et al.*, *NDS 148*, 1 (2018)

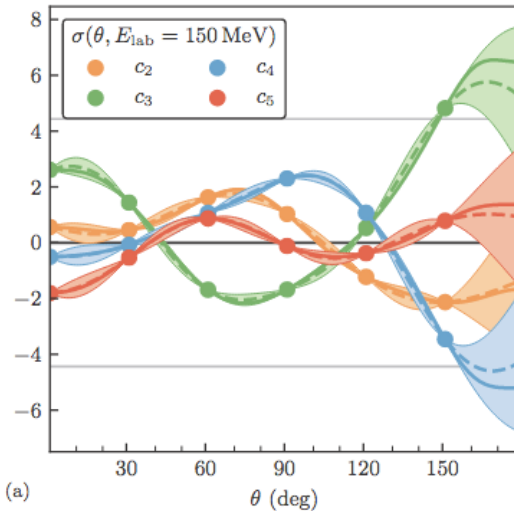
Unc. source	Typical range	Cor(Exp _i ,Exp _i)	Cor(Exp _i , Exp _j) $i \neq j$
$\delta N_{(a)}$	> 1%	Full	$\neq 0$ if same technique/sample
$\delta N_{(b\&c)}$	0-0.5% (Vapor-deposited target) 1% (Painted/electro-plated target)	Full Full	$\neq 0$ if same technique/sample $\neq 0$ if same technique/sample
δc	Eqs. (3) and (5)	Diagonal	0
$\delta\beta$ & δm ; δm	0.02-2%	Gaussian [20]	0.5-0.75
$\delta\beta$ & δm ; $\delta\beta$	0.2-1%	Gaussian	0.5-0.75
$\delta\varepsilon$ & $\delta\alpha$; $\delta\varepsilon$	1.1-4%	Close to full	0.5-1
$\delta\varepsilon$ & $\delta\alpha$; $\delta\alpha$	Compare to nuclear data	Gaussian	0.75-1.0
δb	0.2->10%	Gaussian	Possible
δE	1%, 1-3 ns (TOF, for given TOF length)	From conversion	Technique-dependent
$\delta\phi$	0%, >1%	0.5-Full	Technique-dependent
$\delta\zeta$	See Table III	0.9-1	0.5-0.75
δd	>0.1%	Full	0



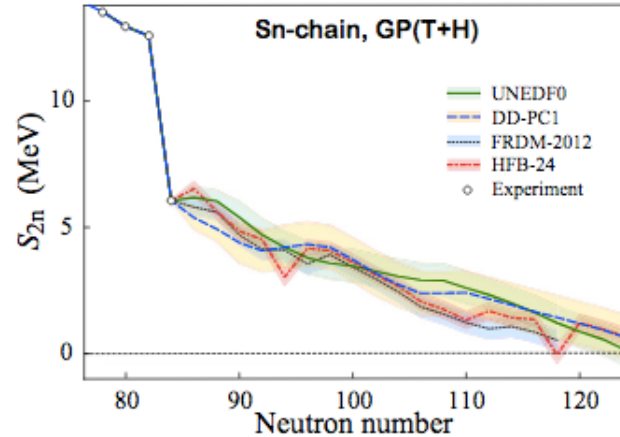
D. Neudecker, *et al.*, *NDS 163*, 228 (2020)

Many groups are exploring machine learning as a method for uncertainty quantification

Gaussian Processes



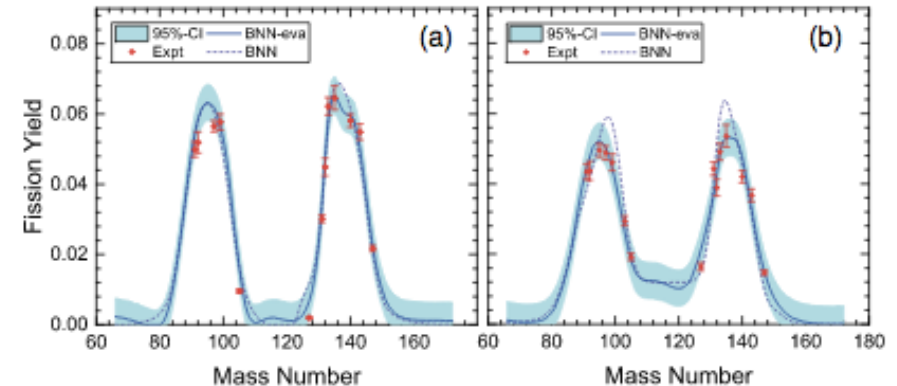
J.A. Melendez, *et al.*, *PRC* **100**, 044001 (2019)



L. Neufcourt, *et al.*, *PRC* **98**, 034318 (2018)

Bayesian Neural Networks

Zi-Ao Wang, *et al.*, *PRL* **123**, 122501 (2019)



Mixture density networks: probabilistic ML algorithm

Input \rightarrow output

$$y = f(x)$$

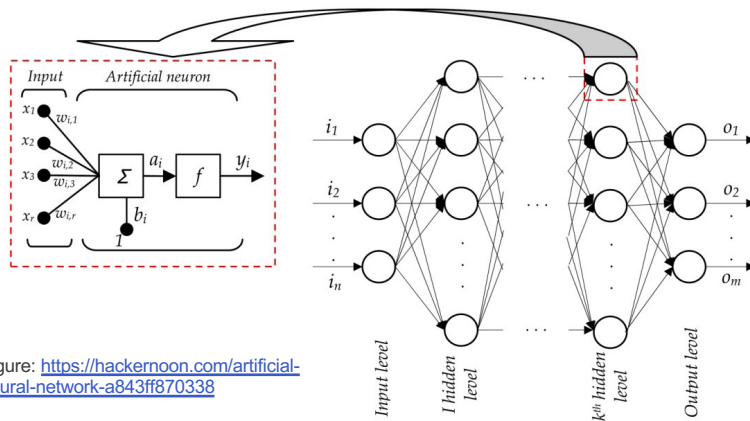
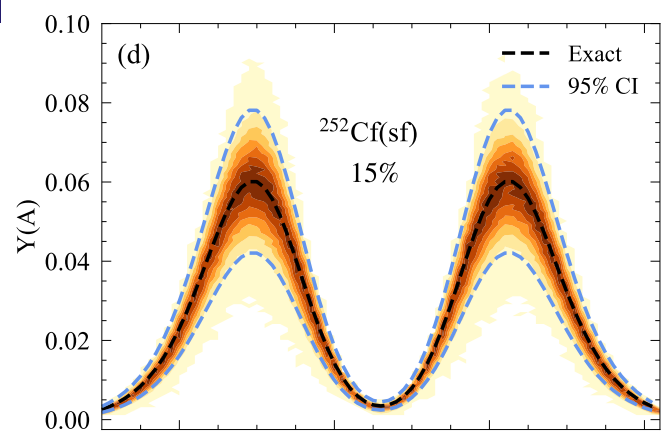


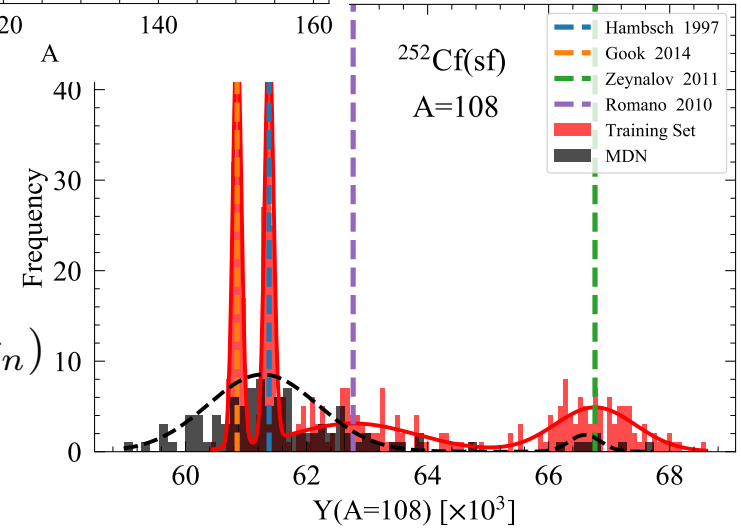
Figure: <https://hackernoon.com/artificial-neural-network-a843ff870338>



A.E. Lovell, A.T. Mohan, P. Talou, *in prep*

$$f(\mathbf{x}) = \alpha_1 \mathcal{N}(\mu_1, \sigma_1) + \alpha_2 \mathcal{N}(\mu_2, \sigma_2) + \dots + \alpha_n \mathcal{N}(\mu_n, \sigma_n)$$

C.M. Bishop, Neural Computing Research Group Report NCRG/94/004 (1994)




Summary

- Uncertainty quantification is gaining more prominence in nuclear theory
- Methods have moved from χ^2 minimization and covariance propagation to Bayesian methods (and are now including machine learning methods as well)
- We have studied both χ^2 minimization and Bayesian methods to constrain uncertainties in the optical potential and found significant differences between the two, particularly in the robustness of the uncertainties and correlations between optical model parameters
- Ongoing studies are beginning to quantify uncertainties from other pieces of the few-body problem

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April 23

Day 8	TBD
9 ₀₀ 10 ₀₀	Exploring experimental conditions to reduce uncertainties in the optical potential EN  Manuel Catacora-Rios Michigan State University