

Systematic nuclear potential for stable, weakly bound and exotic nuclei reactions

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The International Research Collaborators



- Laboratorio Tandar - Comisión Nacional de Energía Atómica (CNEA)

<u>Brazil</u>:

ERSIDAD

 Laboratorio Aberto de Física Nuclear (LAFN) – Universidade de São Paulo⁴⁶ (IFUSP)

Italy:

- Laboratori Nazionale del Sud (LNS) Istituto Nazionale di Fisica Nucleare (INFN).
- Università degli Studi di Padova

<u>Spain</u>:

- Universidad de Sevilla (US)
- Universidad de Huelva (UHU)

<u>Portugal:</u>

- Faculty of Sciences of Lisbon (LIP)

<u>Mexico:</u>

- Universidad Nacional Autónoma de México (UNAM) <u>Costa Rica:</u>
- Universidad de Costa Rica CICANUM





















- **1. The Laboratories and SETUP's**
- 2. The Theoretical Optical Model
- **3. Experimental Data x Theoretical Calculations**
- 4. Conclusions



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Local Research Team:

- A. Di Pietro
- P. Figuera
- R. Spartá
- D. Torresi
- M. La Cognata

<u>Theoretical Support</u>:

- **University of Seville (US) (M. R. Gallardo)**
- Università degli Studi di Padova (J.Casal)
- □ University of Seville (US) (A.M. Moro)
- **D** INFN-Pisa (Jin Lei)

14MV pelletron-tandem accelerator



 $-{}^{7}Li+{}^{119}Sn$ (2017) $-{}^{10}Be+{}^{120}Sn$ (03/2020)



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20MV pelletron-tandem accelerator



Permanent dedicated setup Long term experimental campaigns

> $-{}^{9}\text{Be}+{}^{120}\text{Sn}$ (2018) $-{}^{9}\text{Be}+{}^{197}\text{Au}$ (2018) $-^{10}B+^{197}Au$ (2019) $-^{12,13}C+^{197}Au$ (2020)

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- Andrés Arazi
- Daniel Hojman
- M. A. Carmona
- **Ezequiel Cárdenas**
- **Guillermo Martín**
- José Fernández Nielo

Theoretical Support:

 \succ

University of Seville (US) (M. R. Gallardo) Università degli Studi di Padova (J. Casal)

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de Energía Atómica

20MV pelletron-tandem accelerator



- ELASTIC SCATTERING
 INELASTIC EXCITATION
 NUCLEON TRANSFER
 L. Gasques *et al.*PRC101 (4), 044604 (2020)
- **FUSION**
- M. Aversa et al., PRC101 (4), 044601 (2020)

Permanent dedicated setup Long term experimental campaigns

> $-{}^{9}Be+{}^{120}Sn$ (2018) $-{}^{9}Be+{}^{197}Au$ (2018) $-{}^{10}B+{}^{197}Au$ (2019) $-{}^{12,13}C+{}^{197}Au$ (2020)

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1 ¹⁰B+¹⁹⁷Au @ 38, 40, 42, 43, 44, 45, 46, 47, 48, 49, 50, 51, 52, 53, 54, 55, 56, 57, 58, 59, 61 MeV



20MV pelletron-tandem accelerator

¹³C+¹⁹⁷Au @ 60, 65 and 70 MeV (2018)



Permanent dedicated setup Long term experimental campaigns $-{}^{9}\text{Be}+{}^{120}\text{Sn}$ (2018) $-{}^{9}\text{Be}+{}^{197}\text{Au}$ (2018) $-^{10}B+^{197}Au$ (2019) $-^{12,13}C+^{197}Au$ (2020) Local Research Team: **Daniel Abriola** Andrés Arazi Daniel Hojman M. A. Carmona **Ezequiel Cárdenas Guillermo Martín** José Fernández Nielo

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FROM AN EXPERIMENTAL POINT OF VIEW

We have four main TOOLS:

- **1.** Access to the facilities.
- 2. High beam intensities.
- **3.** Permanent/dedicated experimental setups.
- 4. Long term experimental campaigns.

RECENT GOOD <u>EXPERIMENTAL DATA</u> ARE AVAILABLE



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GOALS:

- ✓ How the weakly (stable or exotic) bound structures affect the dynamics of nuclei reactions?
- ✓ May theoretical models developed to study <u>stable nuclei</u> be applied to <u>weakly bound & exotic</u> ones?

CONCEPTS:

- □ The elastic scattering is the simplest process to test a theoretical model.
- **The optical model (OM)** is the most used theoretical approach for the corresponding **DATA** analysis.

 $U_N(R,E)pprox U_{opt}(R)=V(R)+iW(R)$





Optical Model

The imaginary part of the nuclear potential represents the absorption of the elastic channel.



$$U_N(R,E)pprox U_{opt}(R)=V(R)+iW(R)$$



$$-rac{\hbar^2}{2\mu}
abla^2\Psi(R)+U(R,E)\Psi(R)=E\Psi(R)$$



$$egin{aligned} &-rac{\hbar^2}{2\mu}
abla^2\Psi(R)+U(R,E)\Psi(R)=E\Psi(R)\ &U(R,E)pprox U_N(R)+V_{Coul}(R)pprox U_{opt(R)}+V_{Coul}(R) \end{aligned}$$



$$-\frac{\hbar^{2}}{2\mu}\nabla^{2}\Psi(R) + U(R, E)\Psi(R) = E\Psi(R)$$

$$U(R, E) \approx U_{N}(R) + V_{Coul}(R) \approx U_{opt(R)} + V_{Coul}(R)$$

$$U_{opt}(R) = V_{bare}(R) + V_{pol}(R) + iW_{pol}(R)$$

$$U_{opt}(R) = V_{bare}(R) + V_{pol}(R) + N_{i}V_{SPP}(R)$$
real imaginary potential imaginary



$$U_{opt}(R) = N_R V_{SPP}(R) + N_i V_{SPP}(R)$$

$$V_{ ext{SPP}}\left(R
ight)=V_{ ext{Fold}}\left(R
ight)e^{-4v^{2}/c^{2}}$$



Target

Zero-range interaction:

The NN interaction range is negligible compared to nuclear density

$$egin{aligned}
u_{NN} &= V_0 \delta(ec{R} - ec{r_1} + ec{r_2}) \ \mathrm{V_0} = & -456 \; \mathrm{MeV} \, \mathrm{fm}^3 \; \mathrm{[1]} \end{aligned}$$

[1] L. C. Chamon et al. Phys.Rev. C66, 014610 (2002).



[1]

$$U_{opt}(R) = N_R V_{SPP}(R) + N_i V_{SPP}(R)$$

$$V_{SPP}(R) = V_{Fold}(R) e^{-4v^2/c^2}$$

$$V_{Fold}(R) = \int \int \rho_1(\vec{r}_1) \rho_2(\vec{r}_2) \nu_{NN}(\vec{R} - \vec{r}_1 + \vec{r}_2) d\vec{r}_1 d\vec{r}_2$$

$$V_{Fold}(\vec{R}) = \int \int \rho_1(\vec{r}_1) \rho_2(\vec{r}_2) \nu_{NN}(\vec{R} - \vec{r}_1 + \vec{r}_2) d\vec{r}_1 d\vec{r}_2$$

$$\frac{R - \vec{r}_1 + \vec{r}_2}{\vec{r}_1 + \vec{r}_2}$$

[1] L. C. Chamon et al. Phys.Rev. C66, 014610 (2002).



Stable nuclei reactions have been successfully described assuming the fundamental double-folding nuclear São Paulo potential (SPP).







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$E_{\mathcal{D}} = E - V_{\mathcal{D}}$	projectile	$V_B({ m MeV})$	$R_B(\mathrm{fm})$	$\hbar w ({ m MeV})$
$\boldsymbol{L}_{Red} = \boldsymbol{L}_{c.m.} \boldsymbol{V}_{B}$				
	$^{4}\mathrm{He}$	14.22	9.48	4.92
	⁶ He	12.78	10.52	3.35
	⁶ Li	19.76	10.16	4.20
	7 Li	19.45	10.34	3.86
	$^{9}\mathrm{Be}$	25.78	10.40	3.93
	$^{10}\mathrm{B}$	32.38	10.34	4.17
	¹⁶ O	50.79	10.56	4.14
	¹⁸ O	50.05	10.74	3.86
<image/>	^o Sn tai	rget		

D Different projectiles $+ {}^{120}$ Sn @ energies below, around and above the respective Coulomb barrier


		4		A 1 A	
		*He	14.22	9.48	4.92
		⁶ He	12.78	10.52	3.35
	F = F = V	⁶ Li	19.76	10.16	4.20
	$E_{Red} - E_{c.m.} - v_B$	7 Li	19.45	10.34	3.86
		${}^{9}\mathrm{Be}$	25.78	10.40	3.93
		$^{10}\mathrm{B}$	32.38	10.34	4.17
		$^{16}\mathrm{O}$	50.79	10.56	4.14
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$U_{opt}(R)$ =	$=V_{SPP}+iW(R)$ $V_{opt}(R)$	$=V_{SPP}+iW$ = $W_0 [1+exp(rac{R-R_o}{a_i}]^{-1}]^{-1}$ 100 MeV $a_i=0.25~fm$	$V_{fus}(R)$ n		

projectile

D Different projectiles $+ {}^{120}$ Sn @ energies below, around and above the respective Coulomb barrier

 $\hbar w (MeV)$

 $R_B(\mathrm{fm})$

 $V_B(MeV)$



$U_{opt}(R)$ =	$= V_{SPP} + iW(R)$	$egin{array}{ll} W_{fus} = W_0 \ W_0 = 100 \; M \end{array}$	$[1+exp(rac{R-R_o}{a_i}]^{-1}$ MeV $a_i=0.25~fn$	n —		
		$U_{opt}(R) =$	$V_{SPP} + iW$	$V_{fus}(R)$	Only Internal A (OIA)	bsorption
			¹⁸ O	50.05	10.74	3.86
			$^{16}\mathrm{O}$	50.79	10.56	4.14
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 $V_B(MeV)$

 $R_B(\mathrm{fm})$

D Different projectiles $+ {}^{120}$ Sn @ energies below, around and above the respective Coulomb barrier

 $\hbar w (MeV)$



$U_{opt}(R) = V_{SPP} + iW(R)$	$W_{fus} = W_0 [\ W_0 = 100 \; M$	(OIA) (OIA) $(eV a_i = 0.25 \ fm$				
	$U_{opt}(R) =$	$V_{SPP}+iW$	$V_{fus}(R)$	Only Internal A	bsorption	
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projectile

 $V_B(MeV)$

 $R_B(\mathrm{fm})$

D Different projectiles $+ {}^{120}$ Sn @ energies below, around and above the respective Coulomb barrier

 $\hbar w ({
m MeV})$



	projectile	$V_B({ m MeV})$	$R_B(\mathrm{fm})$	$\hbar w ({ m MeV})$
	$^{4}\mathrm{He}$	14.22	9.48	4.92
	$^{6}\mathrm{He}$	12.78	10.52	3.35
F_{-} - F_{-}	\mathbf{V}_{-} ⁶ Li	19.76	10.16	4.20
$E_{Red} - E_{c.m}$	$h_{\rm Li} = v_B$ ⁷ Li	19.45	10.34	3.86
	$^{9}\mathrm{Be}$	25.78	10.40	3.93
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	$U_{opt}(R) = V_{SPP} + i$	$W_{fus}(R)$	Only Internal A (OIA)	bsorption
$U_{opt}(R) = V_{SPP} + i W(R)$ -	$egin{aligned} U_{opt}(R) &= V_{SPP} + i \ W_{fus} &= W_0 [1 + exp(rac{R-R_o}{a_i}]] \ W_0 &= 100 \; MeV \; a_i = 0.25 \end{aligned}$	$W_{fus}(R) _{]^{-1}} fm$	Only Internal A (OIA)	Absorption
$U_{opt}(R) = V_{SPP} + i W(R)$ -	$egin{aligned} U_{opt}(R) &= V_{SPP} + i \ W_{fus} &= W_0 [1 + exp(rac{R-R_o}{a_i}]] \ W_0 &= 100 \ MeV \ a_i &= 0.25 \end{aligned}$ $U_{opt}(R) &= V_{SPP} + i \end{aligned}$	$W_{fus}(R)$ $^{]^{-1}}_{fm}$ $^{I}N_iV_{SPP}(R)$	Only Internal A (OIA) Strong Surface (SSA)	Absorption

□ Different projectiles + ¹²⁰Sn @ energies below, around and above the respective Coulomb barrier





M. A. G. Alvarez et al., PHYSICAL REVIEW C 100, 064602 (2019)





M. A. G. Alvarez et al., PHYSICAL REVIEW C 100, 064602 (2019)



Experimental DATA from:





M. A. G. Alvarez et al., PHYSICAL REVIEW C 100, 064602 (2019)

OIA

SSA











M. A. G. Alvarez et al., PHYSICAL REVIEW C 100, 064602 (2019)



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M. A. G. Alvarez et al., PHYSICAL REVIEW C 100, 064602 (2019)





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M. A. G. Alvarez et al., PHYSICAL REVIEW C 100, 064602 (2019)



<u>Strong Surface Absorption</u> <u>(SSA)</u>

$$U_{opt}(R) = V_{SPP} + i N_i V_{SPP}(R)$$

Studying the intensity of the imaginary potential

$$egin{aligned} V_{ ext{SPP}}(R) &= V_{ ext{Fold}}(R) e^{-4v^2/c^2} \ V_{ ext{Fold}}(R) &= \int \int
ho_1(ec{r}_1)
ho_2(ec{r}_2)
u_{NN}(ec{R} - ec{r}_1 + ec{r}_2) \, dec{r}_1 \, dec{r}_2 \end{aligned}$$



<u>Sub-barrier region</u>







<u>Sub-barrier region</u>







M. Aversa et al., Phys. Rev. C 101, 044601 (2020)

FIG. 9. (Color online) Experimental reduced fusion cross section as a function of the reduced energy, for several systems involving the same target nucleus: ¹⁹⁷Au. The solid lines represent the reduced BPM cross section.

Fusion is somehow favored for ¹⁰B, while peripheral reaction channels, which are connected to strong surface absorption processes, are favored for the other weakly bound nuclei.























<u>Sub-barrier region</u>























Above-barrier region








<u>Above-barrier region</u>





	1 0	
system	cluster	$\varepsilon_b ({ m MeV})$
11 Li	⁹ Li+2n	-0.396
$^{11}\mathrm{Be}$	$^{10}Be+1n$	-0.501
⁶ He	⁴ He+2n	-0.975
⁶ Li	α +d	-1.473
⁹ Be	⁸ Be+1n	-1.574
⁹ Be	$lpha{+}lpha{+}1\mathrm{n}$	-1.664
$^{7}\mathrm{Li}$	α +t	-2.467
⁹ Li	⁷ Li+2n	-4.062
$^{10}\mathrm{B}$	$^{6}\text{Li}+\alpha$	-4.461
¹⁸ O	$^{14}\mathrm{C} + \alpha$	-6.228
¹⁶ O	$^{12}\mathrm{C}+lpha$	-7.162
$^{12}\mathrm{C}^{g.s.}$	${}^{8}\mathrm{Be}+lpha$	-7.367
$^{12}\mathrm{C}^{H.s.}$	$^{8}\mathrm{Be}{+}lpha$	-7.653
$^{4}\mathrm{He}$	³ He+1n	-20.6

TABLE I. Break-up threshold of lig	ht stable and exotic nuclei.
------------------------------------	------------------------------

projectile	reaction products	$Q({ m MeV})$
	101	
°Li	121 Sn + α + p	2.472
6Li	$^{121}\mathrm{Sb} + \alpha + n$	2.092
7	100	
⁷ Li	122 Sn + α + p	4.036
⁷ Li	$^{122}\mathrm{Sb} + \alpha + n$	1.247
⁹ Be	121 Sn + α + α	4 597
→ ⁹ Be	120 Sn + 8 Be + n	4.505
10 D	121 Cm $+$ 2 cm $+$ m	1 020
D 10 D	$\operatorname{Sn} + 2\alpha + p$	-1.989
1°B	121 Sb + $2\alpha + n$	-2.368



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Λ_{10} B	121 Sn $\pm 2\alpha \pm n$	-1 080
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	, <u> </u>	



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°L1 61.	$121 \operatorname{Sn} + \alpha + p$	2.472
<u> </u>	121 Sb + α + n	2.092
7 T ;	122 C m \perp α \perp m	4 026
Δ1 ⁷ Τ ;	$\sin + \alpha + p$ $122 \operatorname{Sb} + \alpha + n$	4.030 1 947
	bb + a + h	1.24!
⁹ Be	121 Sn + α + α	4.597
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10 B	121 Sp $\pm 2\alpha \pm m$	1 080
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7 Li	$lpha{+}{ m t}$	-2.467
⁹ Li	⁷ Li+2n	-4.062
$^{10}\mathrm{B}$	6 Li+ α	-4.461
¹⁸ O	$^{14}\mathrm{C}+lpha$	-6.228
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$^{12}\mathrm{C}^{g.s.}$	${}^{8}\mathrm{Be}{+}lpha$	-7.367
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system	cluster	$\varepsilon_b ~({ m MeV})$
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11 Be	¹⁰ Be+1n	-0.501
⁶ He	⁴ He+2n	-0.975
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A simple analytical formula for the Coulomb dipole
polarization (CDP) potential was derived in:
M. V. Andrés et al., Nucl. Phys. A 579, 273 (1994).
M. V. Andrés et al., Nucl. Phys. A 583, 817 (1995).

$$U_{pol} = -\frac{4\pi}{9} \frac{Z_t^2}{\hbar v} \frac{1}{(r-a_0)^2 r} \int_{\varepsilon_b}^{\infty} d\varepsilon \frac{dB(E1,\varepsilon)}{d\varepsilon} \left(g\left(\frac{r}{a_0}-1,\xi\right)+if\left(\frac{r}{a_0}-1,\xi\right)\right)$$

 $\frac{dB(E1)}{d\varepsilon}$ is the probability distribution of the dipolar electric transition; $\varepsilon_{\rm b}$ is the necessary energy to break up the projectile; a_0 is the half of the distance of closest approach in the head - on collision; f,g are analytic functions.



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Experimental and theoretical B(E1) distributions



T. Nakamura et al., Phys. Rev. Lett. 96, 252502 (2006). N. Fukuda et al., Phys. Rev. C 70, 054606 (2004).



 $U_{opt}(R) = V_{SPP}(R) + iN_I V_{SPP}(R) + V_{pol}(R) + iW_{pol}(R)$ Coulomb dipole polarization potential



$$U_{opt}(R) = V_{SPP}(R) + iN_I V_{SPP}(R) + V_{pol}(R) + iW_{pol}(R)$$
 Coulomb dipole polarization potential

$$U_{opt}(R) = V_{SPP}(R) + i0.78V_{SPP}(R)$$



 $V_{ ext{SPP}}\left(R
ight)=V_{ ext{Fold}}\left(R
ight)e^{-4v^{2}/c^{2}}$

 $V_{
m Fold}\left(R
ight) = \int \int
ho_{1}(ec{r}_{1})
ho_{2}(ec{r}_{2})
u_{NN}(ec{R}-ec{r}_{1}+ec{r}_{2})\,dec{r}_{1}\,dec{r}_{2}$



$$U_{opt}(R) = V_{SPP}(R) + iN_{I}V_{SPP}(R) + V_{pol}(R) + iW_{pol}(R)$$
Coulomb dipole polarization potential
$$U_{opt}(R) = V_{SPP}(R) + i0.78V_{SPP}(R) \qquad U_{pol} = -\frac{4\pi}{9}\frac{Z_{\iota}^{2}}{\hbar^{V}}\frac{1}{(r-a_{0})^{2}r}\int_{\varepsilon_{b}}^{\infty} d\varepsilon \frac{dB(E1,\varepsilon)}{d\varepsilon} \left(g\left(\frac{r}{a_{0}}-1,\xi\right)+if\left(\frac{r}{a_{0}}-1,\xi\right)\right)$$
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$$\int nuclear potential$$

$$V_{SPP}(R) = V_{Fold}(R)e^{-4v^{2}/c^{2}}$$

$$V_{Fold}(R) = \int \int \rho_{1}(\vec{r}_{1})\rho_{2}(\vec{r}_{2})\nu_{NN}(\vec{R} - \vec{r}_{1} + \vec{r}_{2}) d\vec{r}_{1} d\vec{r}_{2}$$

E (MeV)



studying the intensity of the imaginary CDP potential

cluster ε_b (MeV) system 0 ^{11}Li 9 Li+2n-0.396¹⁰Be+1n $^{11}\mathrm{Be}$ -0.501 -0,1 ⁶He ⁴He+2n -0.975⁶Li $\alpha + d$ -1.473 ⁹Be 8 Be+1n /_{Pol}(MeV) -1.574 -0,2 ⁹Be -1.664 $\alpha + \alpha + 1n$ $- {}^{9}\text{Be} + {}^{208}\text{Pb}; E_{\text{LAB}} = 40.0 \text{ MeV}$ ⁷Li -2.467 $\alpha + t$ \geq $- {}^{6}\text{He} + {}^{120}\text{Sn}; E_{\text{LAB}} = 18.0 \text{ MeV}$ ⁹Li 7 Li+2n -4.062 -0,3 $- {}^{11}\text{Be} + {}^{197}\text{Au}; E_{\text{LAB}} = 39.6 \text{ MeV}$ ^{10}B $^{6}\text{Li} + \alpha$ -4.461 $^{14}C+\alpha$ ^{18}O -6.228- ¹¹Li+²⁰⁸Pb; E_{LAB} = 29.8 MeV -0,4 $^{12}C+\alpha$ ^{16}O -7.162 $^{12}C^{g.s.}$ $^{8}\text{Be}+\alpha$ -7.367 $^{12}\mathrm{C}^{H.s.}$ $^{8}\text{Be}+\alpha$ -7.653 -0,5 20 22 24 26 28 30 12 14 18 $^{4}\mathrm{He}$ 16 3 He+1n -20.6 R(fm)

 $\ensuremath{\mathsf{TABLE}}\xspace$ I. Break-up threshold of light stable and exotic nuclei.



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system	cluster	$\varepsilon_{\rm L}$ (MeV)	
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$^{7}\mathrm{Li}$	$lpha{+}{ m t}$	-2.467	$Be^{-9}Be^{-200}Pb; E_{LAB} = 40.0 \text{ MeV}$
⁹ Li	⁷ Li+2n	-4.062	-0.3 -0.3
$^{10}\mathrm{B}$	6 Li+ $lpha$	-4.461	$-\frac{11}{\text{Be}+197}\text{Au; E}_{\text{H}} = 39.6 \text{ MeV}$
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			K (1111)

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- **1. The Laboratories and SETUP's**
- 2. The Theoretical Optical Model
- **3. Experimental Data x Theoretical Calculations**
- 4. Conclusions



CONCLUSIONS:

- □ We work on an international network based on using pelletron-tandem facilities.
- □ We have analysed 40 elastic scattering angular distributions of stable strongly bound, weakly bound and exotic nuclei, mainly on ¹²⁰Sn, at energies below, around and above the Coulomb barriers.
- □ We have analysed them with an OP based on the nuclear São Paulo Potential, allowing the most sensitive parameter to vary, which showed to be strongly connected to the projectile binding energy.
- Within this context, we have also showed the importance of the Coulomb dipole polarization potential, derived from the semi-classical theory of Coulomb excitation.
- □ We identified the evolution of *long-range absorption* as a function of the projectile binding energy, which, for more exotic nuclei, is dominated by the Coulomb interaction.
- □ The proposed approach shows to be a fundamental basis to study any nuclear reaction.



PHYSICAL REVIEW C 100, 064602 (2019)

Systematic study of optical potential strengths in reactions on ¹²⁰Sn involving strongly bound, weakly bound, and exotic nuclei

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THANK YOU FOR YOUR ATTENTION