

Dissecting calculations of breakup and transfer reactions with halo nuclei using Halo EFT

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Preliminary comments

The different references are hyperlinked to the original papers.
When possible, the link to the arXiv page is also provided.

I'd like to thank the organisers for their excellent initiative to set up this series of international seminars : Jin Lei, Mario Gómez Ramos, Kaitlin Cook, and Jesús Casal.

Thanks for joining in for this on-line seminar.

I wish you all to stay healthy and go through this quarantine unscathed.

Pierre Capel

1 Introduction : Reactions with Halo nuclei

2 Description of ^{15}C

- Single-particle description
- EFT description

3 Reactions with ^{15}C

- Transfer
- Coulomb breakup
- Radiative capture

4 Extension to KO

5 Summary

Halo nuclei

Halo nuclei are found far from stability

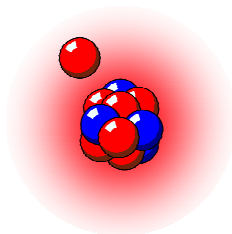
Exhibit peculiar quantal structure :

- Light, **n-rich** nuclei
- Low S_n or S_{2n}

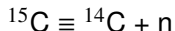
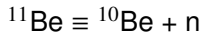
With **large matter radius**

due to strongly clusterised structure :

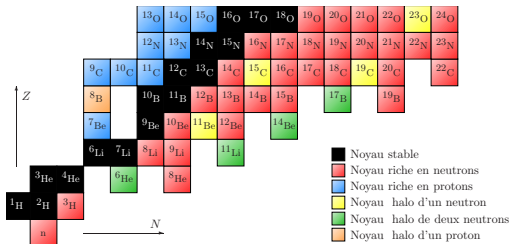
neutrons tunnel far from the **core** and form a **diffuse halo**



One-neutron halo



Two-neutron halo



Two-neutron halo nuclei are **Borromean**. . . [see M. Gómez's , A. Corsi's talks]

This exotic structure challenges nuclear-structure models

Reactions with halo nuclei

Halo nuclei are **fascinating** objects
but difficult to study [$\tau_{1/2}(^{15}\text{C}) = 2.5 \text{ s}$]

How can one **probe their structure** ?

⇒ require **indirect** techniques, like reactions :

- transfer : $^{14}\text{C}(d,p)^{15}\text{C}$
- breakup : $^{15}\text{C} + \text{Pb} \rightarrow ^{14}\text{C} + n + \text{Pb}$
- knockout : $^{15}\text{C} + \text{C} \rightarrow ^{14}\text{C} + \text{X}$
- radiative capture : $^{14}\text{C}(n,\gamma)^{15}\text{C}$ (astrophysical interest)

Need good understanding of the reaction mechanism

(i.e. a good **reaction model**)

to know what nuclear-structure **information** is probed

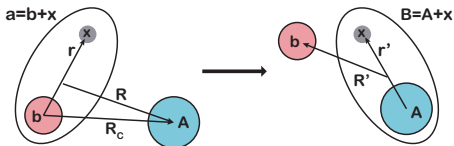
We address this by coupling precise reaction models with **Halo EFT**

^{11}Be : [P.C., Phillips, Hammer PRC 98, 034610 (2018)] arXiv

^{15}C : [Moschini, Yang, P.C. PRC 100, 044615 (2019)] arXiv

Few-body description of transfer

Transfer : $A + a(\equiv b + x) \rightarrow B(\equiv A + x) + b$ aka $A(a, x)B$
 described in a **few-body** model :



Bound states, initial Φ_{bx} and final Φ_{Ax}
 obtained from effective interactions : V_{bx} and V_{Ax}

Scattering described by optical potentials U_{aA} , U_{pB}, \dots

For (d, p)

$$T_{\text{post}}(p B, d A) = \langle \chi_{pB}^{(-)} \Phi_{An} | V_{pn} + U_{pA} - U_{pB} | \Psi_{dA}^{(+)} \rangle$$

@FR-ADWA $\Psi_{dA}^{(+)} \approx \chi_{dA}^{(+)} \Phi_{pn}$ [Johnson & Tandy NPA 235, 56 (1974)]
 [see also M. Gómez Ramos', A. Ratkiewicz's and N. Timofeyuk's talks]

Few-body description of breakup

Breakup : $B(\equiv A + n) + T \rightarrow A + n + T$

Projectile (B) \equiv **core** (A)+loosely bound **neutron** (n) described by

$$H_{An} = T_r + V_{An}(\mathbf{r})$$

V_{An} effective interaction

describes the A - n system

with ground state Φ_{An}

Target T seen as structureless

Interaction with target simulated by optical potentials

\Rightarrow breakup reduces to **three-body** scattering problem :

$$[T_R + H_{An} + U_{AT} + U_{nT}] \Psi(\mathbf{r}, \mathbf{R}) = E_T \Psi(\mathbf{r}, \mathbf{R})$$

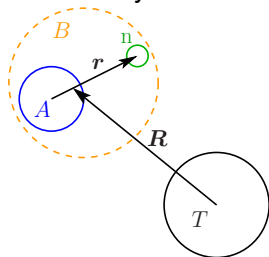
with initial condition $\Psi(\mathbf{r}, \mathbf{R}) \xrightarrow{Z \rightarrow -\infty} e^{iKZ} \Phi_{An}(\mathbf{r})$

We use the Dynamical Eikonal Approximation (DEA) @ 70AMeV

[Baye, P. C., Goldstein, PRL 95, 082502 (2005)]

An eikonal approximation with **relativistic corrections** @ 600AMeV

[Moschini, P. C. PLB 790, 367 (2019)] arXiv



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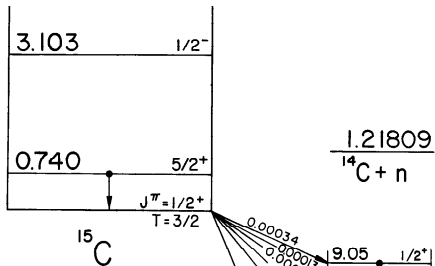
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$$^{15}\text{C} \equiv ^{14}\text{C} \otimes n$$



- $\frac{1}{2}^+$ ground state :
 $\epsilon_{\frac{1}{2}^+} = -1.218 \text{ MeV}$
 Seen as $1s_{1/2}$ neutron bound to $^{14}\text{C}(0^+)$
- $\frac{5}{2}^+$ bound excited state :
 $\epsilon_{\frac{5}{2}^+} = -0.478 \text{ MeV}$
 Seen as $0d_{5/2}$ neutron bound to $^{14}\text{C}(0^+)$

Single-particle description

In reaction models, one-neutron halo nucleus \equiv **two-body** system :

$$H_{An} = T_r + V_{An}(\mathbf{r}),$$

where V_{An} is a phenomenological Woods-Saxon that reproduces the basic nuclear properties of the projectile (binding energy, J^π, \dots)

The halo-nucleus structure is described by H_{An} eigenstates

$$H_{An} \Phi_{nljm}(\mathbf{r}) = \epsilon_{nlj} \Phi_{nljm}(\mathbf{r}),$$

$$\Phi_{nljm}(\mathbf{r}) = \frac{1}{r} u_{nlj}(r) \mathcal{Y}_{ljm}(\Omega)$$

Asymptotically, $u_{nlj}(r) \xrightarrow{r \rightarrow \infty} C_{nlj} e^{-\kappa_{nlj} r}$ with $\hbar^2 \kappa_{nlj}^2 / 2\mu_{An} = |\epsilon_{nlj}|$

C_{nlj} is the **Asymptotic Normalisation Coefficient** (ANC)

In A -n continuum $u_{klj}(r) \xrightarrow{r \rightarrow \infty} kr \sin(kr + \delta_{lj})$ with $\hbar^2 k^2 / 2\mu_{An} = \epsilon$

δ_{lj} is the **phaseshift** in partial wave lj

Halo-EFT potential

Replace A - n interaction by **effective** potentials in each partial wave

Use **Halo EFT** : clear separation of scales (in energy or in distance)

⇒ provides an expansion parameter (small scale / large scale)

along which the low-energy behaviour is expanded

[C. Bertulani, H.-W. Hammer, U. Van Kolck, NPA 712, 37 (2002)] arXiv

[H.-W. Hammer, C. Ji, D. R. Phillips JPG 44, 103002 (2017)] arXiv

Use narrow Gaussian potentials @ LO@ NLO

$$V_{lj}(r) = V_0^{lj} e^{-\frac{r^2}{2r_0^2}} + V_2^{lj} r^2 e^{-\frac{r^2}{2r_0^2}}$$

r_0 used to evaluate the sensitivity of calculations to short-range physics

@ LO :

Fit V_0^{lj} to reproduce ϵ_{nlj} (known experimentally)

@ NLO :

Fit V_2^{lj} to also reproduce C_{nlj} (obtained by analysing transfer reactions)

[Yang, P.C. PRC 98, 054602 (2018)] arXiv

LO analysis of transfer reaction $^{14}\text{C}(d,p)^{15}\text{C}$

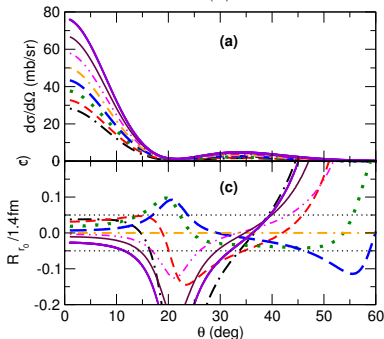
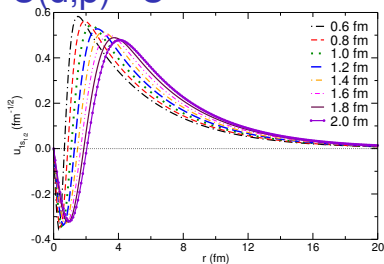
Transfer is purely **peripheral**
 (\Leftrightarrow probe only **ANC**) at

- low E_d
- forward angles

[Yang, P.C. PRC 98, 054602 (2018)] arXiv

- 1 Generate different wave functions @ LO using different r_0
- 2 Run transfer calculations (we use FR-ADWA)
- 3 Determine the zone of peripherality

$$R_{r_0/1.4\text{fm}} = \left(\frac{C_{1s1/2}^{(1.4\text{fm})}}{C_{1s1/2}^{(r_0)}} \right)^2 \frac{d\sigma^{(r_0)}/d\Omega}{d\sigma^{(1.4\text{fm})}/d\Omega} - 1$$



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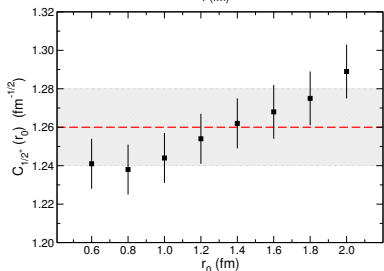
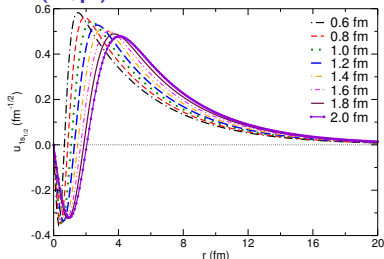
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- 4 Within that zone, fit to data and extract an **ANC** :

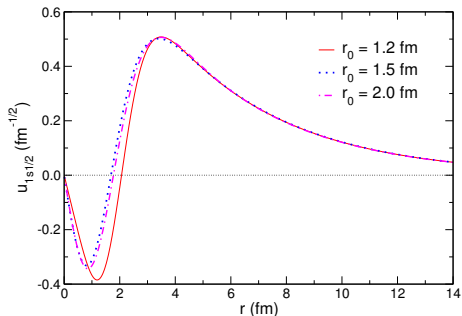
$C_{1s1/2} = 1.26 \pm 0.02 \text{ fm}^{-1/2}$, close to *ab initio* $C_{1s1/2} = 1.28 \text{ fm}^{-1/2}$



$s_{\frac{1}{2}}$: @ NLO potentials fitted to $\epsilon_{\frac{1}{2}^+}$ and $C_{1s1/2}$

Potentials fitted to $\epsilon_{1s1/2} = -1.218$ MeV and $C_{1s1/2} = 1.26$ fm $^{-1/2}$

Ground-state wave function



Th : [Moschini, Yang, P.C. PRC 100, 044615 (2019)] arXiv

- Wave functions : **same** asymptotics but **different** interior
- In higher partial waves ($l \geq p$) $V_{lj} = 0$

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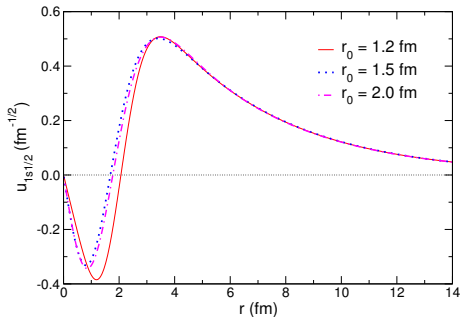
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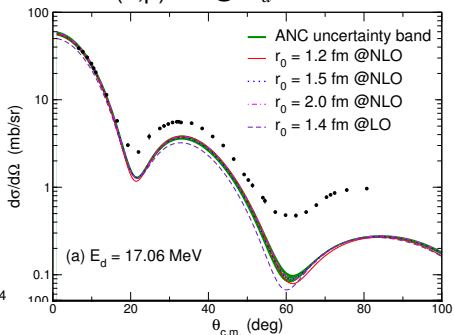
$^{14}\text{C}(d,p)^{15}\text{C}$ @ $E_d = 17\text{MeV}$

Potentials fitted to $\epsilon_{1s1/2} = -1.218\text{ MeV}$ and $C_{1s1/2} = 1.26\text{ fm}^{-1/2}$

Ground-state wave function



$^{14}\text{C}(d,p)^{15}\text{C}$ @ $E_d = 17\text{MeV}$

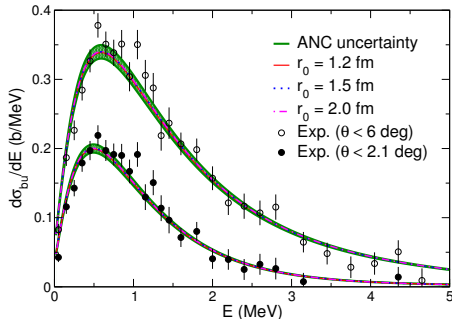


Th : [Moschini *et al.* PRC 100, 044615 ('19)] arXiv Exp : [Mukhamedzhanov PRC 84, 024616 (2011)]

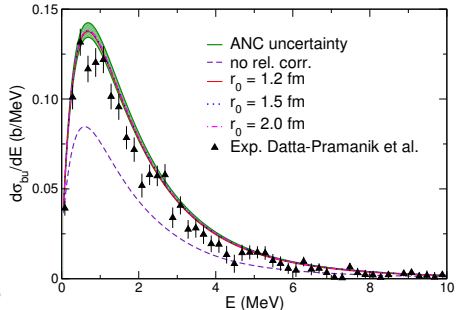
- Wave functions : **same** asymptotics but **different** interior
- **Excellent agreement** with low energy data at forward angle, where little sensitivity to r_0

Coulomb breakup : $^{15}\text{C} + \text{Pb} \rightarrow ^{14}\text{C} + n + \text{Pb}$

RIKEN : 68A MeV

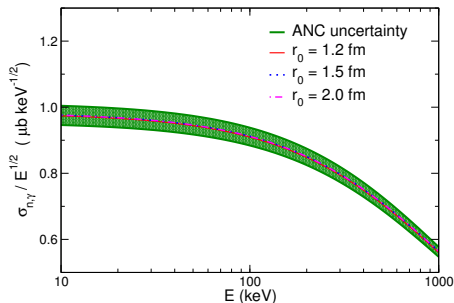
Exp : [Nakamura *et al.* PRC 79, 035805 (2009)]DEA : [Baye *et al.* PRL 95, 082502 (2005)]

GSI : 605A MeV

Exp : [Datta Pramanik *et al.* PLB 551, 63 (2003)]

Eikonal : [Moschini, P. C. PLB 790, 367 (2019)]

- All calculations provide **very similar** results for all r_0 despite the difference in the internal part of the wave function \Rightarrow reaction is **peripheral** [P.C. & Nunes PRC75, 054609 (2007)] arXiv
- **Excellent** agreement with data confirms **ANC** and $\delta_p \approx 0$ [Moschini *et al.* PRC 100, 044615 ('19)] arXiv

Radiative capture $^{14}\text{C}(n,\gamma)^{15}\text{C}$ 

E (keV)	$\sigma_{n,\gamma}^{\text{exp}}$ (μb)	$\sigma_{n,\gamma}^{\text{th}}$ (μb)
23.3	7.1 ± 0.5	5.8 ± 0.2
150	10.7 ± 1.2	10.6 ± 0.3
500	17.0 ± 1.5	15.4 ± 0.4
800	15.8 ± 1.6	16.7 ± 0.5

Exp : [Reifarth *et al.* PRC 77, 015804 (2008)]

Th : [Moschini *et al.* PRC 100, 044615 ('19)] arXiv

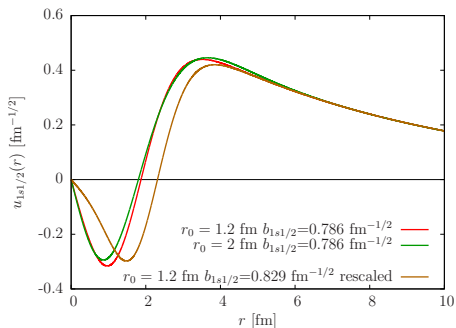
- Insensitive to r_0 , because purely peripheral
- $\sigma_{n,\gamma}$ in agreement with *ab initio* prediction
- Good **agreement** with the data but at low energy...

Extension of the idea to KO

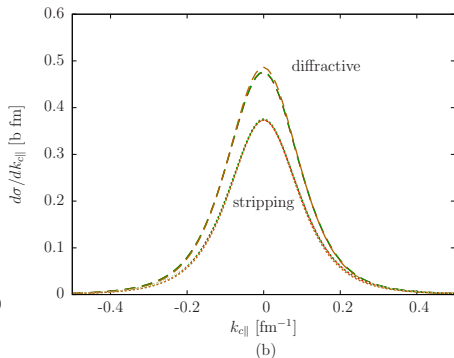
Theoretical analysis of KO : $^{11}\text{Be} + \text{C} \rightarrow ^{10}\text{Be} + \text{X}$ @ 70 A MeV

[Hebborn & P.C. PRC 100, 054607 (2019)] arXiv

Halo-EFT wave functions



KO cross section



- Using NLO **Halo-EFT** description of ^{11}Be within eikonal model of KO
- **Same** $d\sigma/dp_{c||}$ with **same ANC** \Rightarrow reaction purely peripheral
- **Insensitive** to description of continuum \Rightarrow good probe of ANC

Summary and prospect

- Halo nuclei studied mostly through reactions
- Mechanism of reactions with halo nuclei understood
How to reliably infer structure information on halos from reactions ?
Halo EFT : [P.C., Phillips, Hammer PRC 98, 034610 (2018)] arXiv
Efficient way to include the significant degrees of freedom : ANC, δ_{lj}
- We apply this idea for ^{15}C [Moschini *et al.* PRC 100, 044615 ('19)] arXiv
Using one Halo-EFT description of ^{15}C , we reproduce
 - ▶ transfer $^{14}\text{C}(d,p)^{15}\text{C}$
 - ▶ Coulomb breakup (intermediate and high energy)
 - ▶ Radiative capture $^{14}\text{C}(n,\gamma)^{15}\text{C}$
- KO is also purely peripheral [Hebborn, P.C. PRC 100, 054607 ('19)] arXiv
- Future :
 - ▶ Analyse existing data on KO
 - ▶ Extend to other nuclei (^{31}Ne)
 - ▶ Include core excitation in Halo EFT

Thanks to my collaborators

Hans-Werner Hammer
Achim Schwenk



Daniel Phillips



Laura Moschini
Jiecheng Yang
Chloë Hebborn

